

A reciprocity relationship for random diffuse fields  
in acoustics and structures with application to the  
analysis of uncertain systems

Robin Langley

## Example 1: Acoustic loads on satellites

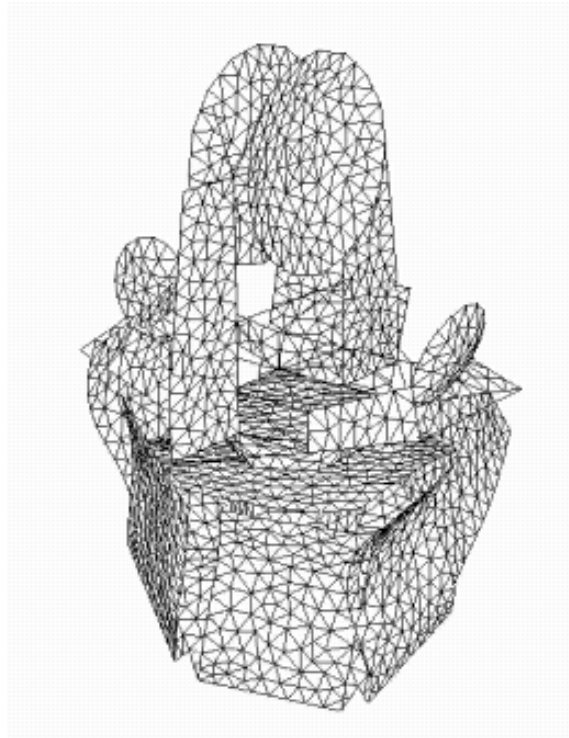


Arianne 5 Launch

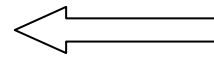


Acoustic test facility at ESTEC

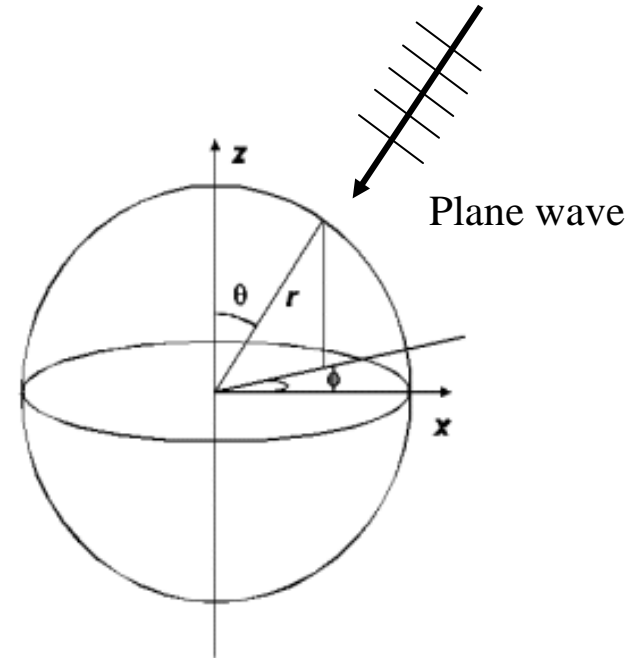
# Satellite computational model



Structural model



Pressure  $p$



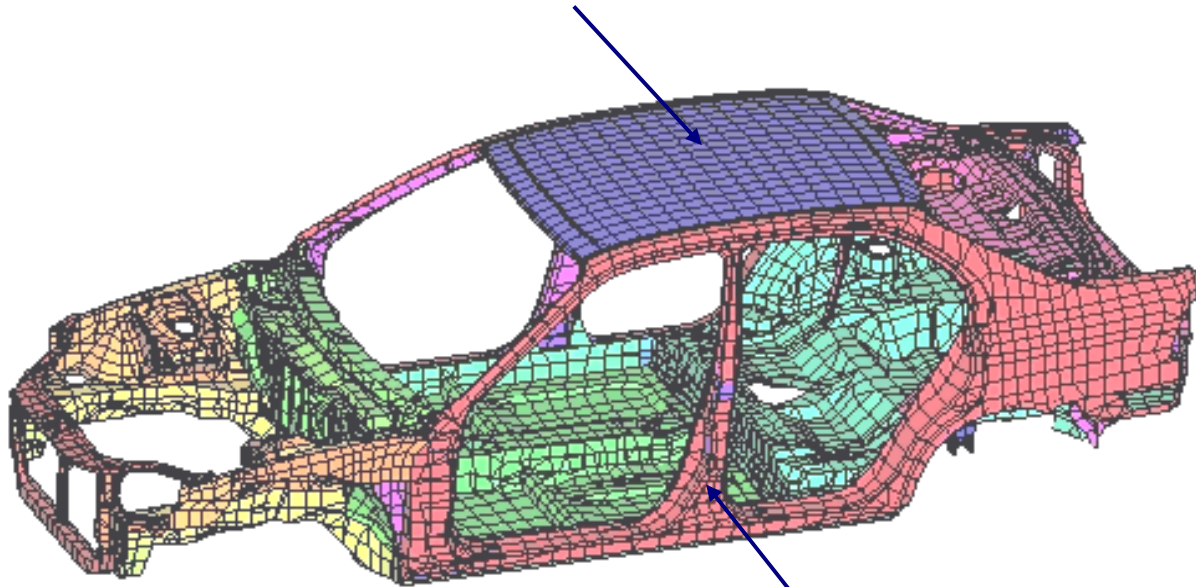
Diffuse acoustic field

$$p = \sum_j \sum_k \alpha_{jk} \exp \{ ikr \sin \theta_j \cos \phi_k + ikr \sin \theta_j \sin \phi_k + ikr \cos \theta_j + i\omega t \}$$

Pressure modelled as a summation of plane waves at each frequency  
This is very time consuming - is there a more efficient approach?

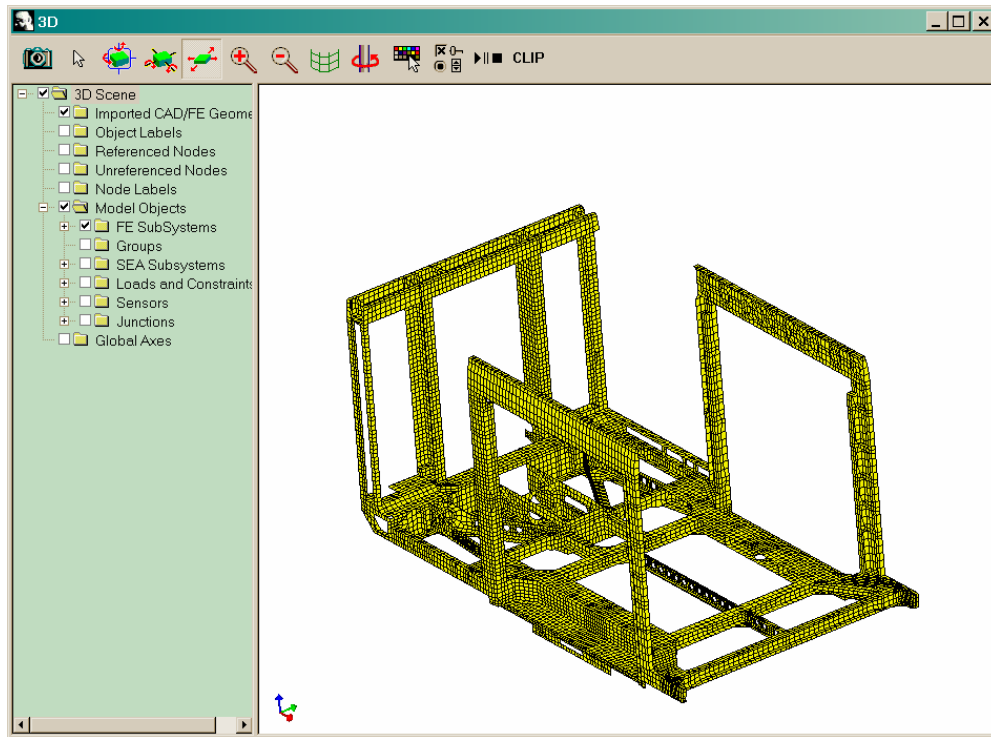
## Example 2: High frequency structural vibrations

Thin panels – very many modes, response can be viewed as a diffuse wave field

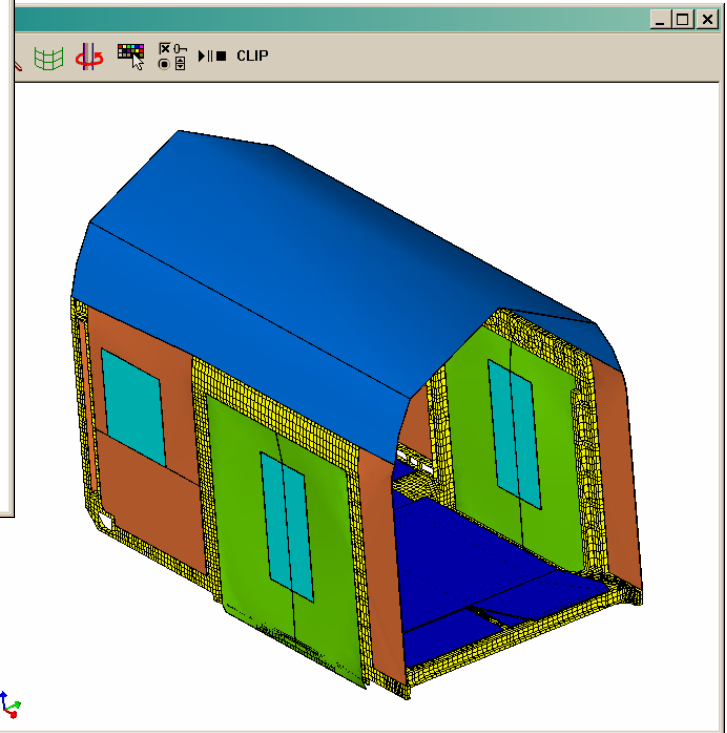


Stiff components – detailed FE model needed

# Structure computational model



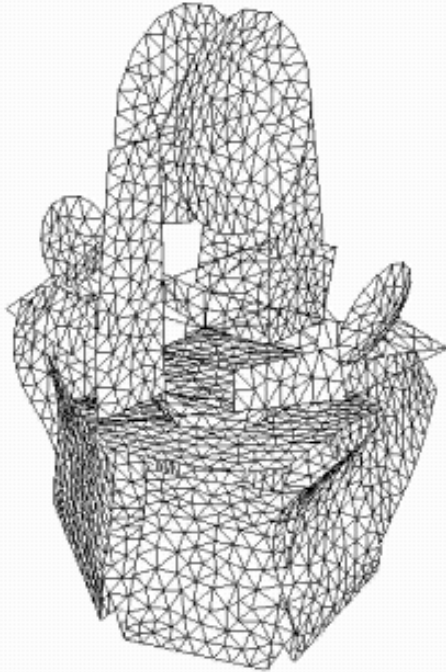
FE model – displacement dofs



SEA model – diffuse wave fields

How do we couple these very different types of model?  
Again we need to compute the forces arising from the diffuse waves.

# The diffuse field reciprocity relation



Finite element model has dof  $\mathbf{x}$

Pressure  $p$  produces force vector  $\mathbf{f}$

Can we avoid summing waves to compute  $\mathbf{f}$  ?

The *diffuse field reciprocity relation* states:

$$\mathbf{S}_{ff} = \mathbf{E} \left[ \mathbf{f} \mathbf{f}^{*T} \right] = \left( \frac{4E}{\omega \pi n} \right) \text{Im} \left\{ \mathbf{D}_{\infty} \right\}$$

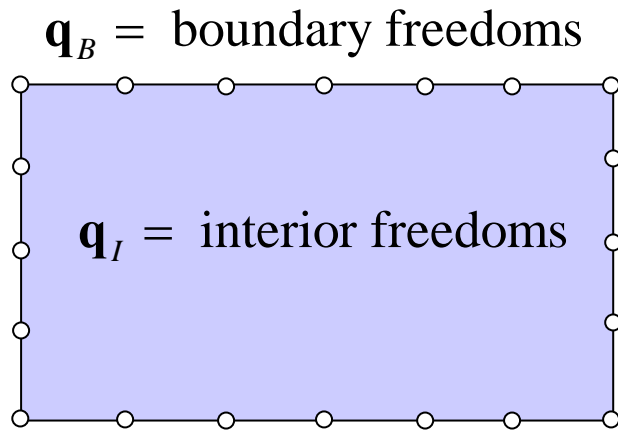
$E$  is the energy of the diffuse field

$n$  is the modal density of the diffuse field

(for acoustics  $E/n$  is specified by the noise level)

$\mathbf{D}_{\infty}$  is the dynamic stiffness matrix for radiation into an *infinite* acoustic system – readily computable

## Short proof of the relation (structural component)



$$\begin{pmatrix} \mathbf{D}_{11} & \mathbf{D}_{12} \\ \mathbf{D}_{12}^T & \mathbf{D}_{22} \end{pmatrix} \begin{pmatrix} \mathbf{q}_B \\ \mathbf{q}_I \end{pmatrix} = \begin{pmatrix} \mathbf{g} \\ \mathbf{0} \end{pmatrix}$$

Express  $\mathbf{q}_I$  in terms of blocked modes

$$\mathbf{q}_I = \Phi \mathbf{a}$$

where  $\mathbf{a}$  are the modal amplitudes

$$\begin{pmatrix} \mathbf{D}_{11} & \mathbf{C}_{12} \\ \mathbf{C}_{12}^T & \Lambda \end{pmatrix} \begin{pmatrix} \mathbf{q}_B \\ \mathbf{a} \end{pmatrix} = \begin{pmatrix} \mathbf{g} \\ \mathbf{0} \end{pmatrix}$$

$$\Lambda_{kk} = \omega_k^2 (1 + i\eta) - \omega^2$$

**Boundary dynamic stiffness**  $\{\mathbf{D}_{11} - \mathbf{C}_{12} \Lambda^{-1} \mathbf{C}_{12}^T\} \mathbf{q}_B = \mathbf{g}$  or  $\mathbf{D} \mathbf{q}_B = \mathbf{g}$

**Force on boundary due to interior waves**

$$\mathbf{f} = -\mathbf{C}_{12} \mathbf{a} \quad \Rightarrow \quad \mathbf{f} \mathbf{f}^{*T} = \mathbf{C}_{12} \mathbf{a} \mathbf{a}^{*T} \mathbf{C}_{12}^T$$

# Ensemble averaging

Assume that:

- the natural frequencies are random
- The modal amplitudes  $\mathbf{a}$  correspond to a diffuse field, i.e. they are statistically independent and each mode stores equal energy

$$\mathbf{D} = \mathbf{D}_{11} - \mathbf{C}_{12} \mathbf{\Lambda}^{-1} \mathbf{C}_{12}^T \Rightarrow \mathbf{E}[\text{Im}(\mathbf{D})] = \mathbf{E}[\mathbf{C}_{12} \mathbf{C}_{12}^{*T}] (\pi n / 2N\omega)$$

$$\mathbf{f} \mathbf{f}^{*T} = \mathbf{C}_{12} \mathbf{a} \mathbf{a}^{*T} \mathbf{C}_{12}^T \Rightarrow \mathbf{E}[\mathbf{f} \mathbf{f}^{*T}] = \mathbf{E}[\mathbf{C}_{12} \mathbf{C}_{12}^T] (2E / N\omega^2)$$

$$\text{Thus: } \mathbf{E}[\mathbf{f} \mathbf{f}^{*T}] = \left( \frac{4E}{\omega\pi n} \right) \mathbf{E}[\text{Im}\{\mathbf{D}\}]$$

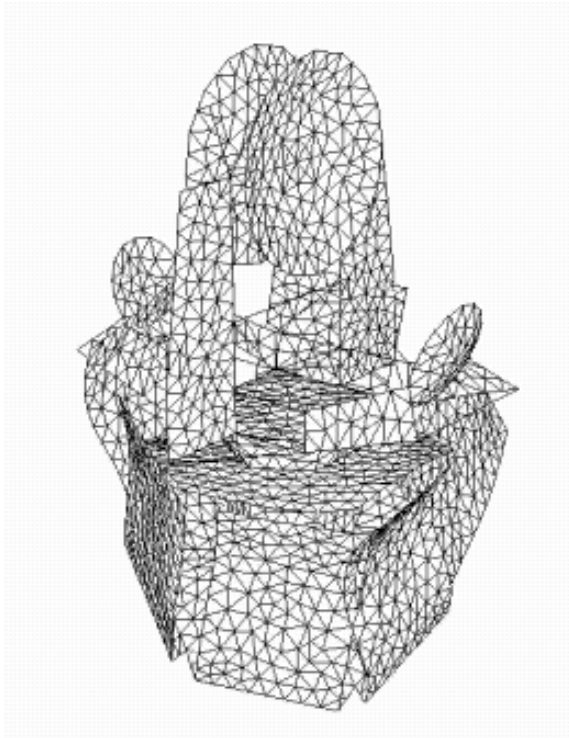
The proof is completed by noting that  $\mathbf{E}[\mathbf{D}] = \mathbf{D}_\infty$

Smith, P. JASA 34, 640-647, 1962

Shorter, P.J. and Langley, R.S. JASA 117, 85-95, 2005



# Application 1: diffuse acoustic loading



$\mathbf{D}_\infty$  is the *radiation* dynamic stiffness of the structure. This can be found by using the *boundary element method*.

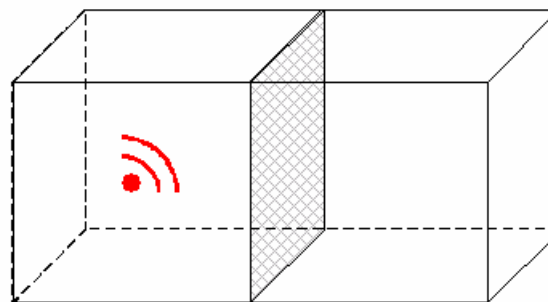
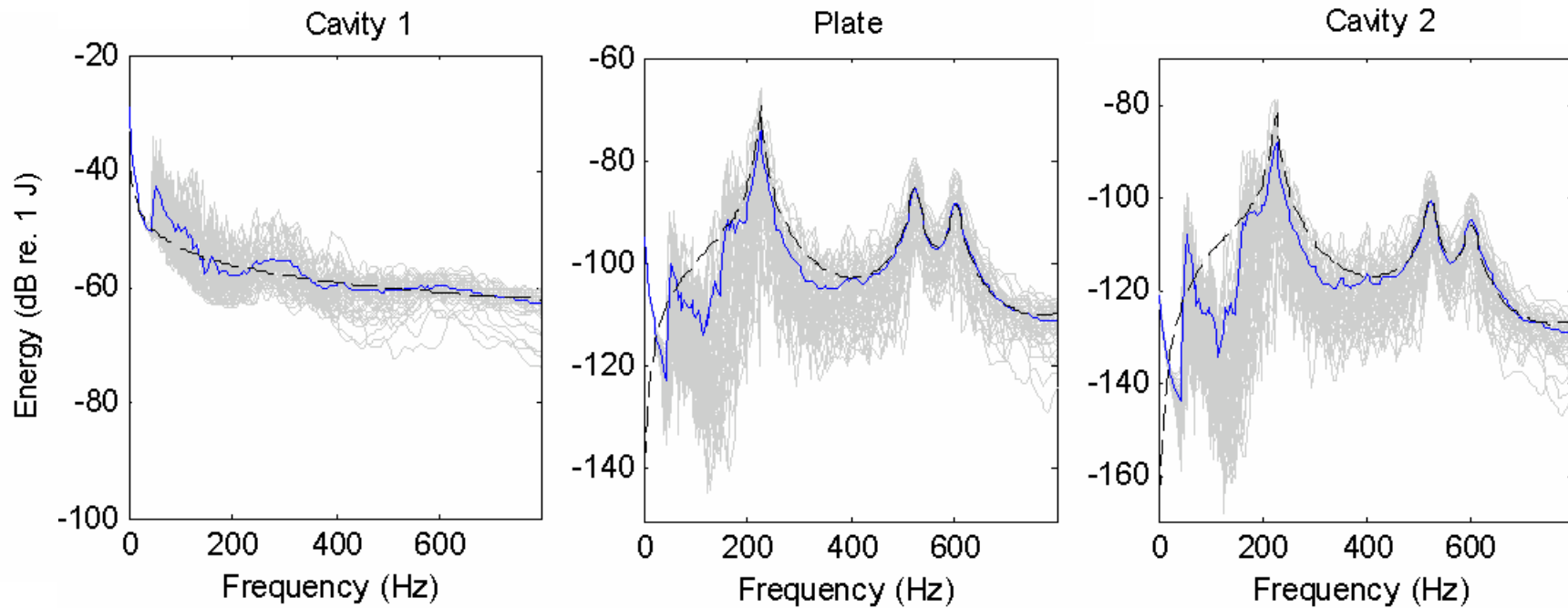
The diffuse field reciprocity relation

$$\mathbf{E}[\mathbf{ff}^{*\text{T}}] = \left( \frac{4E}{\omega\pi n} \right) \mathbf{E}[\text{Im}\{\mathbf{D}\}]$$

replaces the need for a direct summation

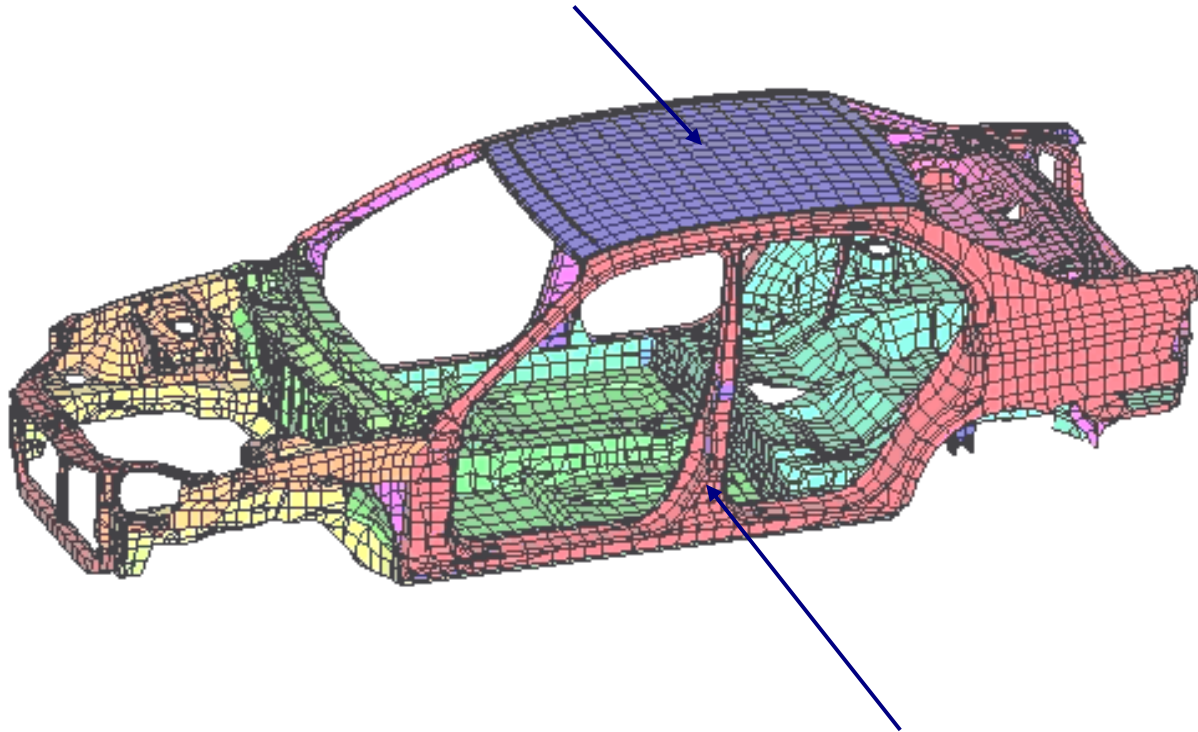
$$p = \sum_j \sum_k \alpha_{jk} \exp\{ikr \sin \theta_j \cos \phi_k + ikr \sin \theta_j \sin \phi_k + ikr \cos \theta_j + i\omega t\}$$

# Example – acoustic transmission, ensemble of random cavities



## Application 2: the hybrid FE-SEA method

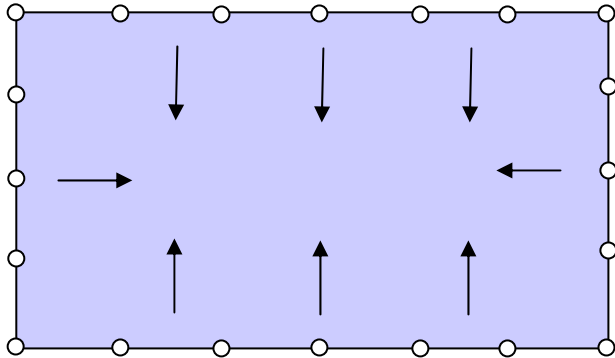
Thin panels – very many modes, response can be viewed as a diffuse wave field. SEA degree of freedom  $E_k$



Stiff components – detailed FE model needed  
FE degrees of freedom  $\mathbf{q}$

# Hybrid Equations – the response of subsystem $k$

## Direct Field

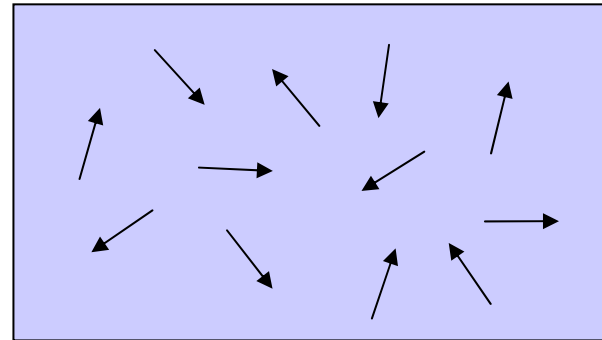


Boundary freedoms  $\mathbf{q}$

Boundary motions generate a “direct field” of ingoing waves. Associated boundary forces written as:

$$-\mathbf{D}_{dir}^{(k)} \mathbf{q}$$

## Reverberant Field



Reflections produce the “reverberant field”. Associated boundary forces written as:

$$\mathbf{f}_{rev}^{(k)}$$

## Hybrid Equations continued

So the system equations are written as:

$$\mathbf{D}_{tot} \mathbf{q} = \mathbf{f} + \sum_k \mathbf{f}_{rev}^{(k)}, \quad \mathbf{D}_{tot} = \mathbf{D}_d + \sum_k \mathbf{D}_{dir}^{(k)}$$

$\mathbf{D}_{dir}^{(k)}$  can be found by considering wave radiation into a semi-infinite system

$\mathbf{f}_{rev}^{(k)}$  is accounted for by using a *diffuse field reciprocity* result:

$$\mathbf{S}_{ff}^{(k),rev} \equiv \mathbf{E} \left[ \mathbf{f}_{rev}^{(k)} \mathbf{f}_{rev}^{(k)*T} \right] = \left( \frac{4E_k}{\omega\pi n_k} \right) \text{Im} \left\{ \mathbf{D}_{dir}^{(k)} \right\}$$

## Final form of the Hybrid Equations

$$\omega(\eta_j + \eta_{d,j})E_j + \sum_k \omega\eta_{jk}n_j(E_j/n_j - E_k/n_k) = P_{in,j}^{ext}$$

$$\mathbf{S}_{qq} = \mathbf{D}_{tot}^{-1} \left[ \mathbf{S}_{ff} + \sum_k \left( \frac{4E_k}{\omega\pi n_k} \right) \text{Im} \left\{ \mathbf{D}_{dir}^{(k)} \right\} \right] \mathbf{D}_{tot}^{-1*T}$$

where:

$$P_{in,j}^{ext} = (\omega/2) \sum_{rs} \text{Im} \left\{ D_{dir,rs}^{(j)} \right\} \left( \mathbf{D}_{tot}^{-1} \mathbf{S}_{ff} \mathbf{D}_{tot}^{-1*T} \right)_{rs}$$

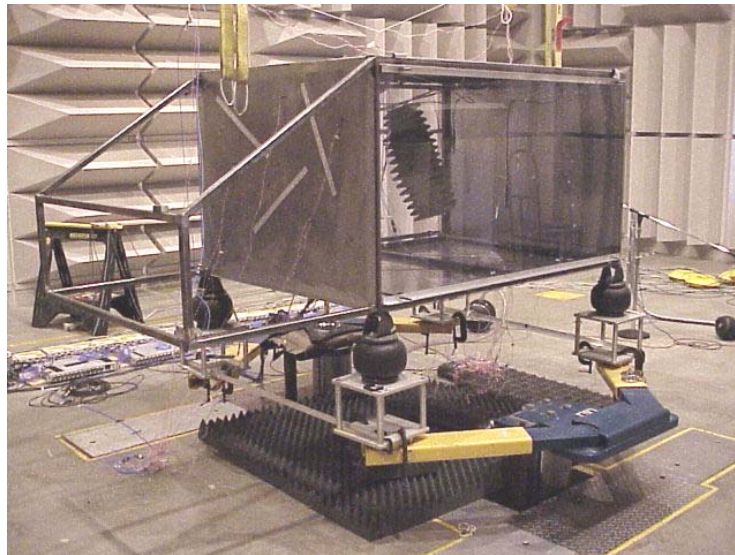
$$\omega\eta_{jk}n_j = (2/\pi) \sum_{rs} \text{Im} \left\{ D_{dir,rs}^{(j)} \right\} \left( \mathbf{D}_{tot}^{-1} \text{Im} \left\{ \mathbf{D}_{dir}^{(k)} \right\} \mathbf{D}_{tot}^{-1*T} \right)_{rs}$$

$$\omega\eta_{d,j} = \left( \frac{2}{\pi n_j} \right) \sum_{rs} \text{Im} \left\{ D_{d,rs} \right\} \left( \mathbf{D}_{tot}^{-1} \text{Im} \left\{ \mathbf{D}_{dir}^{(j)} \right\} \mathbf{D}_{tot}^{-1*T} \right)_{rs}$$

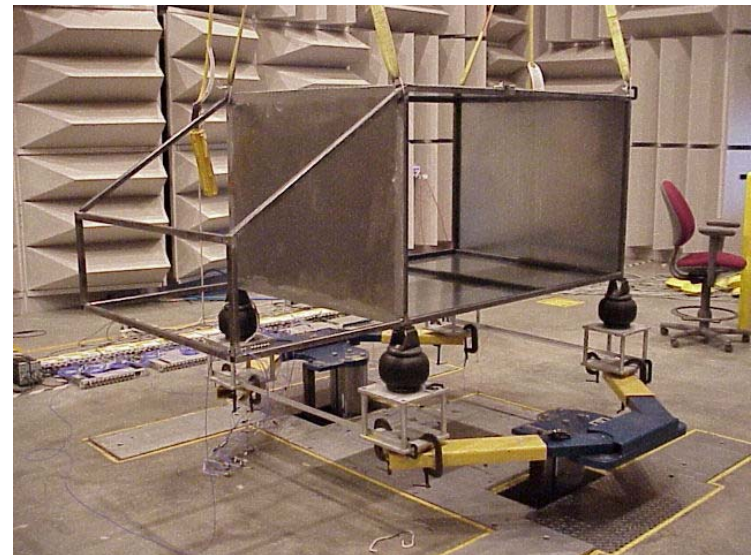
# Example application of the hybrid method

## Test box

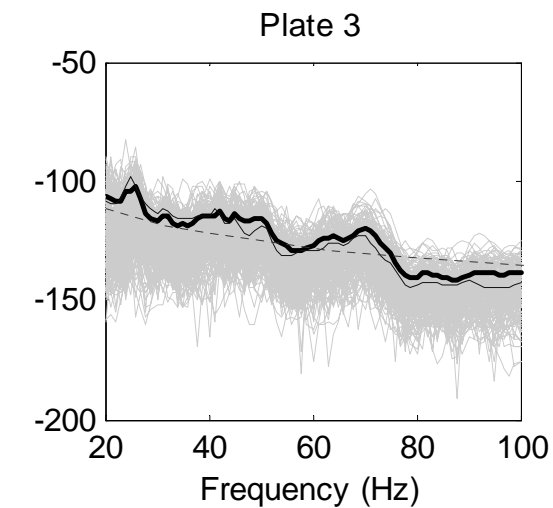
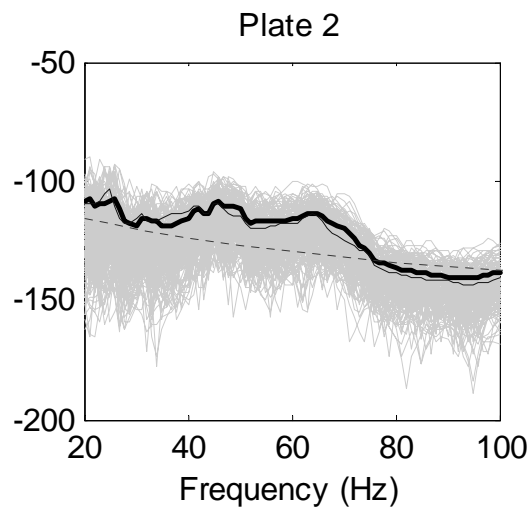
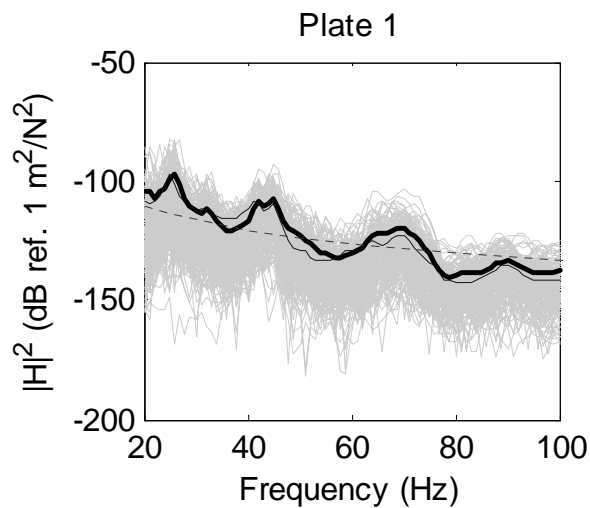
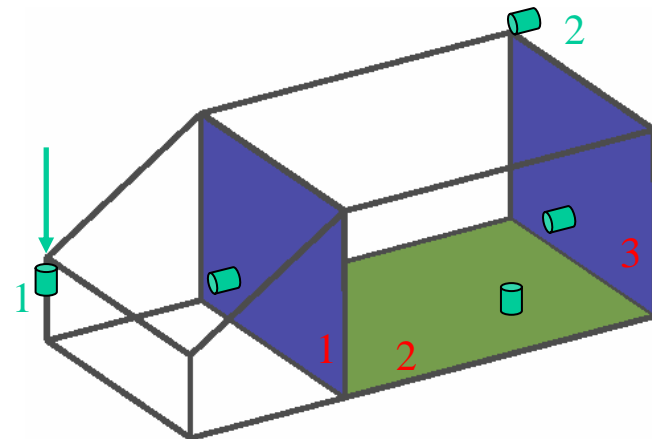
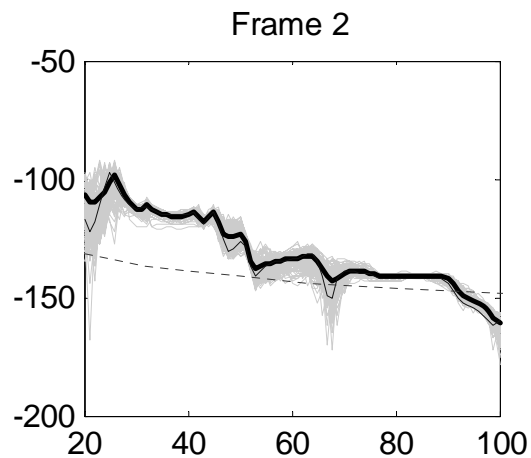
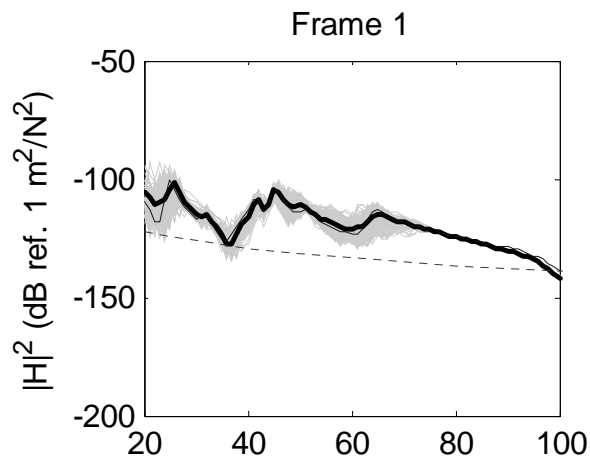
- frame: welded beam sections 1"x1"x1/8"
- 5 plates of different thickness



Courtesy of General Motors Corporation



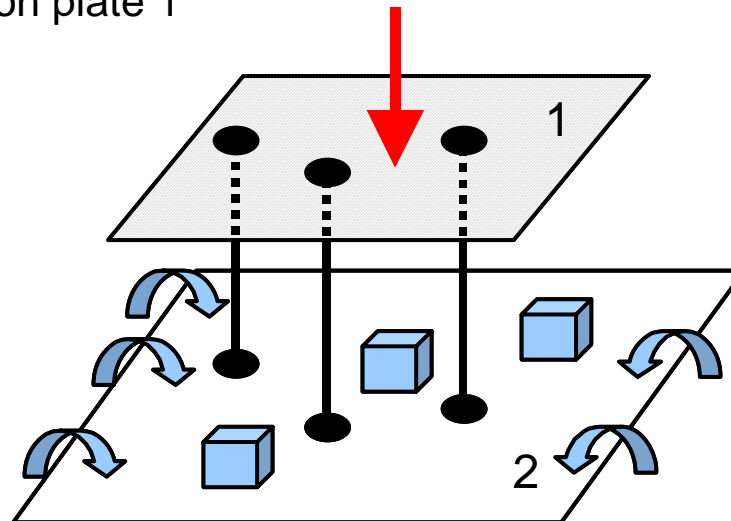
# Example application of the hybrid FE-SEA method



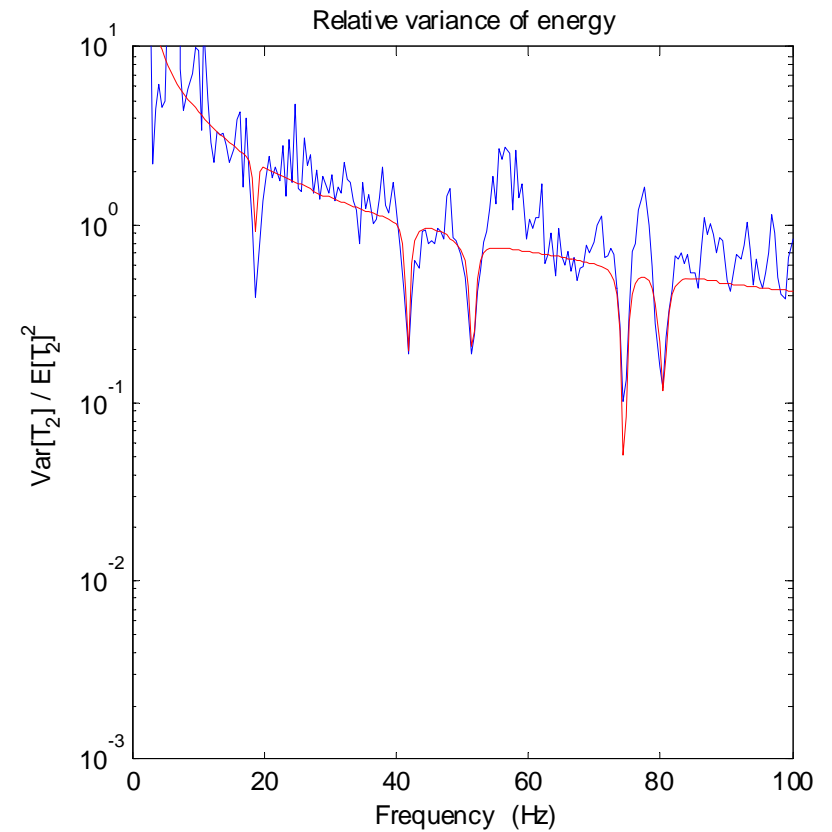
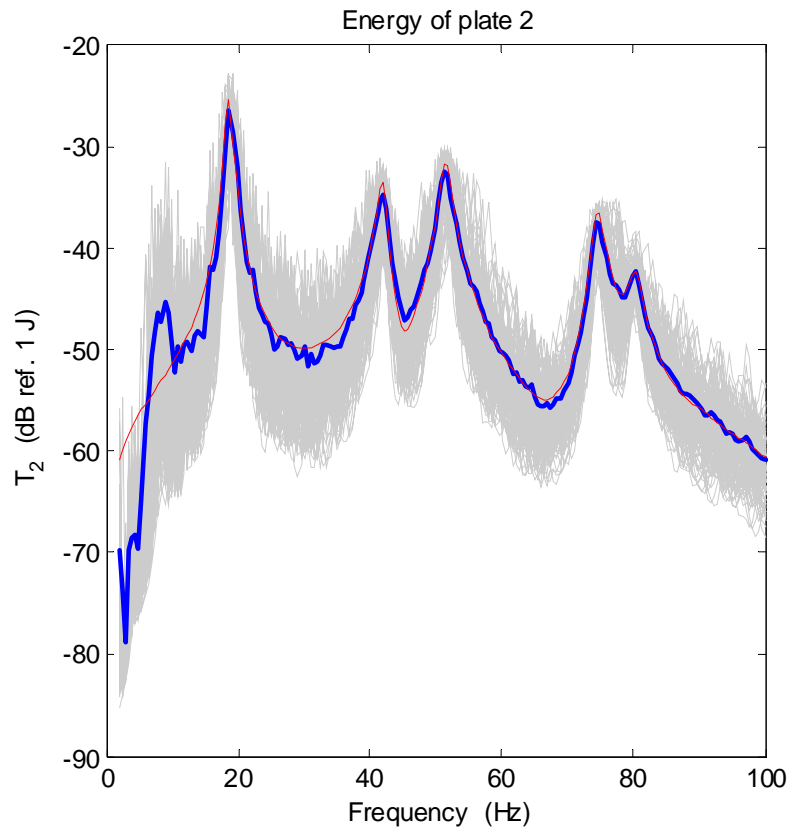


# Further Development: Variance Prediction in the Hybrid Method

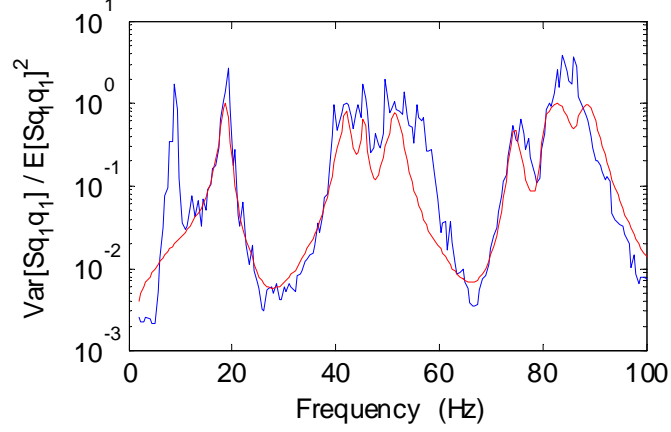
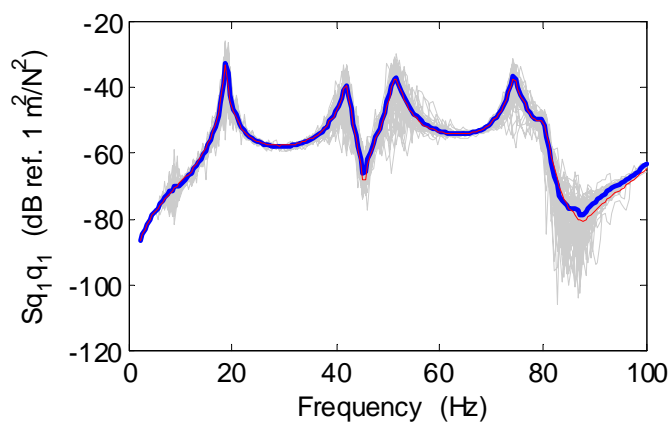
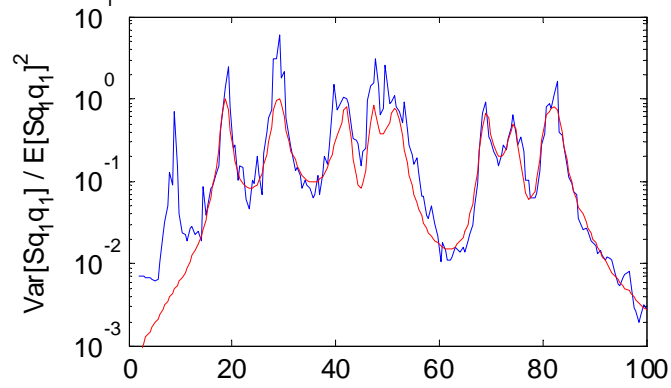
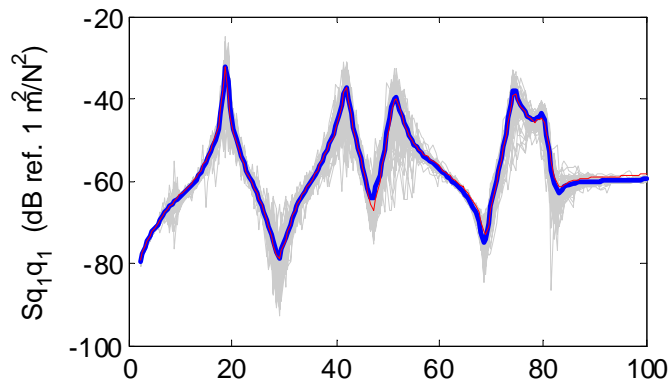
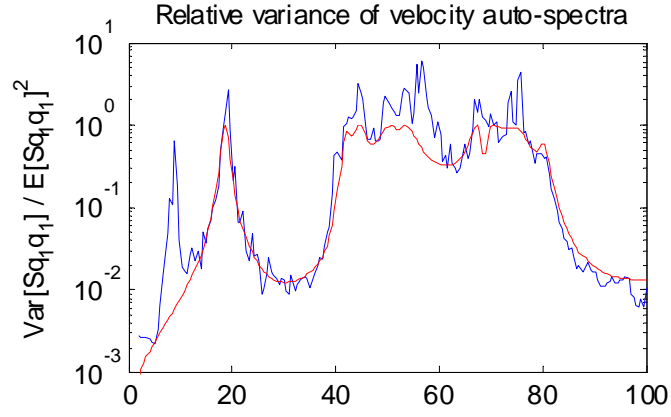
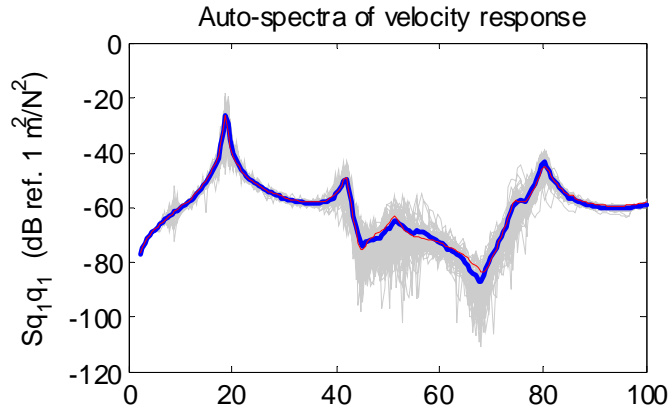
- Research example: two plates coupled through 3 points
- Plate 1 is deterministic  
(long wavelength,  $\sim 9$  modes below 100 Hz)
- Plate 2 is randomized by masses and edge springs  
(short wavelength,  $\sim 148$  modes below 100 Hz)
- Harmonic point force on plate 1



# Ensemble mean & relative variance of plate 2 energy response



Ensemble mean &  
relative variance  
auto-spectra  
response at  
3 points on plate 1



# Postscript

For systems that are driven purely at the boundary, it is unlikely that the excitation will produce a perfect diffuse field. In this case it has recently been found that the reciprocity relationship must be extended:

$$E[\mathbf{ff}^{*T}] = \left( \frac{4E}{\pi\omega n} \right) \text{Im}(\mathbf{D}_\infty) + \left( \frac{2}{\pi m} \right) \left\{ 2 \text{Re}(\mathbf{S}_{\hat{f}\hat{f}}) + q(m)\mathbf{S}_{\hat{f}\hat{f}} \right\}$$

This extension clears up an apparent anomaly between the reciprocity relationship and recent results concerning the variance of the response of a random system.

Langley, R.S. JASA submitted, 2006